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Conformal Lie 2-algebras and conformal omni-Lie algebras

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Abstract. The notions of conformal Lie 2-algebras and conformal omni-Lie algebras are introduced. It is proved that the category of conformal Lie 2-algebras and the category of 2-term conformal L_∞ -algebras are equivalent. We construct conformal Lie 2-algebras from conformal omni-Lie algebras and Leibniz conformal algebras.

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1. Introduction

The algebraic theory of Lie 2-algebras, which constitutes a categorification of Lie algebras, was extensively developed by Baez and Crans [BC04]. As a truncated version of L_∞ -algebras, Lie 2-algebras replace the underlying vector space with a 2-term chain complex of vector spaces (a 2-vector space) and the Jacobi identity with a natural transformation called the Jacobiator, subject to specific coherence laws. Baez and Crans established an equivalence between the category of Lie 2-algebras and that of 2-term L_∞ -algebras. For further developments in this area, we refer to [LL20; LSZ14; SL13; SL16].

Omni-Lie algebras, introduced by Weinstein [Wei00], represent a central example of Lie 2-algebras. They arise as linearizations of Courant algebroids [LWX97], and since every Courant algebroid induces a Lie 2-algebra, omni-Lie algebras naturally fall within this framework. Subsequent research has extended their study to various generalizations, including omni-Lie algebroids and omni-Lie 2-algebras, as explored in [CL10; KW01; SLZ11]. More recently, [ZL14] investigated the Dirac structures of omni-Lie superalgebras.

In this paper, we consider whether a categorification of Lie conformal algebras or vertex Lie algebras exists. We answer this affirmatively by introducing the notion of a conformal Lie 2-algebra, which is a Lie conformal algebra object in the category of $\mathbb{C}[\partial]$ -modules internal to 2-vector spaces. We also define 2-term conformal L_∞ -algebras and prove an equivalence between the category of conformal Lie 2-algebras and that of 2-term conformal L_∞ -algebras.

The second part of the paper is devoted to constructing examples of conformal Lie 2-algebras. In Section 4.3, we introduce conformal omni-Lie algebras and use them to produce conformal Lie 2-algebras. Moreover, in the final part of Section 4, we provide a construction of conformal Lie 2-algebras from Leibniz conformal algebras.

To illustrate this construction, we also provide explicit examples, as not every Leibniz conformal algebra arises from an omni-Lie algebra.

The paper is organized as follows. In Section 2, we recall background material on Lie conformal algebras and Leibniz conformal algebras. Section 3 introduces conformal Lie 2-algebras and 2-term conformal L_∞ -algebras, and establishes the categorical equivalence between them. Section 4 examines special cases of conformal Lie 2-algebras, such as skeletal and strict types. We also define the conformal omni-Lie algebra $\mathcal{E}$ associated with a Lie conformal algebra and a representation. Finally, we construct conformal Lie 2-algebras from conformal omni-Lie algebras and Leibniz conformal algebras.

Throughout this paper, we denote by $\mathbb{C}$ the field of complex numbers. Let M be a vector space. The space of polynomials of λ with coefficients in M is denoted by $M[\lambda]$.

2. Preliminaries

In this section, we will recall some facts and definitions about Lie conformal algebras and conformal Leibniz algebras, see [BKV99; DK98; GT14; HB20; Kac98; LTW13; Lib08; Zha15] for more details.

A *conformal algebra* R is a $\mathbb{C}[\partial]$ -module with a λ -product $\cdot_\lambda \cdot$ which defines a $\mathbb{C}$ -linear map from $A \otimes A \rightarrow A[\lambda]$ satisfying the conformal sesquilinearity condition

$$(\partial x)_\lambda y = -\lambda x_\lambda y, \quad x_\lambda \partial y = (\partial + \lambda)x_\lambda y, \quad (1)$$

for all $x, y \in A$.

Definition 2.1. A *Lie conformal algebra* R is a conformal algebra with the $\mathbb{C}$ -bilinear map $[\cdot_\lambda \cdot] : R \times R \rightarrow R[\lambda]$ satisfying

$$[x_\lambda y] = -[y_{-\partial-\lambda} x], \quad (\text{skew-symmetry}) \quad (2)$$

$$[x_\lambda [y_\mu z]] = [[x_\lambda y]_{\lambda+\mu} z] + [y_\mu [x_\lambda z]], \quad (\text{Jacobi identity}) \quad (3)$$

for all $x, y, z \in R$.

Using the skew-symmetry condition, we can rewrite the Jacobi identity as the following identity:

$$[[x_\lambda y]_{\lambda+\mu} z] = [x_\lambda [y_\mu z]] + [[x_\lambda z]_{-\partial-\mu} y]. \quad (4)$$

Definition 2.2. A *left Leibniz conformal algebra* $(R, [\cdot_\lambda \cdot])$ is a conformal algebra satisfying the following left Leibniz identity

$$x \circ_\lambda (y \circ_\mu z) = (x \circ_\lambda y) \circ_{\lambda+\mu} z + y \circ_\mu (x \circ_\lambda z). \quad (5)$$

Definition 2.3. A *right Leibniz conformal algebra* $(R, [\cdot_\lambda \cdot])$ is a conformal algebra satisfying the following right Leibniz identity

$$(x \circ_\lambda y) \circ_{\lambda+\mu} z = x \circ_\lambda (y \circ_\mu z) + (x \circ_\lambda z) \circ_{-\partial-\mu} y. \quad (6)$$

Remark 2.4. Using the equivalent identity (4), we can see a Lie conformal algebra as both a left and a right Leibniz conformal algebra. However, it should be noted that

the converse is not true as the general Leibniz conformal algebra does not meet the requirement of skew-symmetry.

Definition 2.5. A *left module* M over a Lie conformal algebra R is a $\mathbb{C}[\partial]$ -module endowed with a map $R \times M \rightarrow M[\lambda]$, $(x, v) \mapsto x \triangleright_\lambda v$, satisfying the following axioms:

$$(\partial x) \triangleright_\lambda v = -\lambda x \triangleright_\lambda v, \quad x \triangleright_\lambda (\partial v) = (\partial + \lambda)x \triangleright_\lambda v, \quad (7)$$

$$[x_\lambda y] \triangleright_{\lambda+\mu} v = x \triangleright_\lambda (y \triangleright_\mu v) - y \triangleright_\mu (x \triangleright_\lambda v), \quad (8)$$

for all $x, y \in R, v \in M$. Similarly, a *right module* M over a Lie conformal algebra R is a $\mathbb{C}[\partial]$ -module endowed with a bilinear map $M \times R \rightarrow M[\lambda]$, $(v, x) \mapsto v \triangleleft_\lambda x$, satisfying the following axioms:

$$(\partial v) \triangleleft_\lambda x = -\lambda v \triangleleft_\lambda x, \quad v \triangleleft_\lambda (\partial x) = (\partial + \lambda)v \triangleleft_\lambda x, \quad (9)$$

$$v \triangleleft_\mu [x_\lambda y] = (v \triangleleft_\mu x) \triangleleft_{\lambda+\mu} y - (v \triangleleft_\mu y) \triangleleft_{-\partial-\lambda} x. \quad (10)$$

A right R -module become a left R -module if we define $x \triangleright_\lambda v$ to be $-v \triangleleft_{-\partial-\lambda} x$.

Definition 2.6 ([Zha15]). A *module or representation* M over a Leibniz conformal algebra $(R, [\cdot_\lambda \cdot])$ is a $\mathbb{C}[\partial]$ -module endowed with two $\mathbb{C}$ -bilinear maps $R \times M \rightarrow M[\lambda]$, $(x, v) \mapsto x \triangleright_\lambda v$ and $M \times R \rightarrow M[\lambda]$, $(v, x) \mapsto v \triangleleft_\lambda x$ satisfying the following axioms ($a, b \in R, v \in M$):

$$(\partial x) \triangleright_\lambda v = -\lambda x \triangleright_\lambda v, \quad x \triangleright_\lambda (\partial v) = (\partial + \lambda)x \triangleright_\lambda v, \quad (11)$$

$$(\partial v) \triangleleft_\lambda x = -\lambda v \triangleleft_\lambda x, \quad v \triangleleft_\lambda (\partial x) = (\partial + \lambda)v \triangleleft_\lambda x, \quad (12)$$

$$x \triangleright_\lambda (y \triangleright_\mu v) = (x \circ_\lambda y) \triangleright_{\lambda+\mu} v + y \triangleright_\mu (x \triangleright_\lambda v), \quad (13)$$

$$x \triangleright_\lambda (v \triangleleft_\mu y) = (x \triangleright_\lambda v) \triangleleft_{\lambda+\mu} y + v \triangleleft_\mu (x \circ_\lambda y), \quad (14)$$

$$v \triangleleft_\lambda (x \circ_\mu y) = (x \triangleleft_\lambda x) \triangleleft_{\lambda+\mu} y + x \triangleright_\mu (v \triangleleft_\lambda y). \quad (15)$$

We denote it by $(M, \triangleright_\lambda, \triangleleft_\lambda)$.

Next, we give some examples of Lie conformal algebras.

Example 2.7. Let $\mathfrak{g}$ be a Lie algebra. Let $\text{Cur } \mathfrak{g} := \mathbb{C}[\partial] \otimes \mathfrak{g}$ be the free $\mathbb{C}[\partial]$ -module. Then $\text{Cur } \mathfrak{g}$ is a Lie conformal algebra, called current Lie conformal algebra, with λ -bracket given by:

$$[(f(\partial) \otimes x)_\lambda (g(\partial) \otimes y)] := f(-\lambda)g(\lambda + \partial) \otimes [x, y],$$

for all $f(\partial), g(\partial) \in \mathbb{C}[\partial], x, y \in \mathfrak{g}$.

The Virasoro Lie conformal algebra Vir is the simplest nontrivial example of Lie conformal algebras. It is defined by

$$\text{Vir} = \mathbb{C}[\partial]L, \quad [L_\lambda L] = (\partial + 2\lambda)L.$$

Let M and N be $\mathbb{C}[\partial]$ -modules. A conformal linear map f from M to N is a $\mathbb{C}$ -linear map $f_\lambda : M \rightarrow \mathbb{C}[\lambda] \otimes N$ such that $f_\lambda \partial = (\partial + \lambda)f_\lambda$. The category of $\mathbb{C}[\partial]$ -modules with conformal linear maps as morphisms is denoted by Vect^∂ .

Let $\text{Chom}(M, N)$ denote the set of conformal linear maps from M to N . Then $\text{Chom}(M, N)$ is a $\mathbb{C}[\partial]$ -module via:

$$\partial f_\lambda = -\lambda f_\lambda.$$

The composition $fg : L \rightarrow N$ of conformal linear maps $f : M \rightarrow N$ and $g : L \rightarrow M$ is given by $(f_\lambda g)_{\lambda+\mu} = f_\lambda g_\mu$. If M is a finitely generated $\mathbb{C}[\partial]$ -module, then $\text{Cend}(M) := \text{Chom}(M, M)$ is an associative conformal algebra with respect to the above composition. Thus, $\text{Cend}(M)$ becomes a Lie conformal algebra, which is denoted as $\text{Cend}(M)$, with respect to the following λ -bracket

$$[f_\lambda g]_\mu = f_\lambda g_{\mu-\lambda} - g_{\mu-\lambda} f_\lambda,$$

equivalently,

$$[f_\lambda g] = f_\lambda g - g_{-\partial-\lambda} f.$$

All $\mathbb{C}[\partial]$ -modules in this paper are assumed to be finitely generated.

A homomorphism between two Lie conformal algebras $(R, [\cdot_\lambda \cdot])$ and $(R', [\cdot_\lambda \cdot]')$ is a conformal linear map $f : R \rightarrow R'$ such that

$$f_\lambda \partial(x) = (\partial + \lambda) f_\lambda(x), \quad f([x_\lambda y]) = [f(x)_\lambda f(y)]',$$

for all $x, y \in R$.

An n -cochain ($n \in \mathbb{Z}_{\geq 0}$) of a Lie conformal algebra R with coefficients in a module M is a $\mathbb{C}$ -linear map

$$\begin{aligned} \gamma : \quad R^{\otimes n} &\longrightarrow M[\lambda_1, \dots, \lambda_{n-1}], \\ (x_1, \dots, x_n) &\longmapsto \gamma_{\lambda_1, \dots, \lambda_{n-1}}(x_1, \dots, x_n) = \gamma(x_{1\lambda_1} \dots x_{n-1\lambda_{n-1}} x_n), \end{aligned}$$

where $M[\lambda_1, \dots, \lambda_{n-1}]$ denotes the space of polynomials with coefficients in M , satisfying the following conditions:

(1) Conformal linearity:

$$\gamma_{\lambda_1, \dots, \lambda_{n-1}}(x_1, \dots, \partial x_i, \dots, x_n) = -\lambda_i \gamma_{\lambda_1, \dots, \lambda_{n-1}}(x_1, \dots, x_i, \dots, x_n).$$

(2) Skew-symmetry:

$$\gamma(x_1, \dots, x_{i+1}, \partial(x_i), \dots, x_n) = -\gamma(x_1, \dots, x_i, \partial(x_{i+1}), \dots, x_n).$$

Let $R^{\otimes 0} = \mathbb{C}$ so that a 0-cochain is an element $v \in M$ and $(\delta v)_\lambda x = x_\lambda v$. Define a differential δ of a cochain R by

$$\begin{aligned} &(\delta \gamma)_{\lambda_1, \dots, \lambda_n}(x_1, \dots, x_{n+1}) \\ &= \sum_{i=1}^{n+1} (-1)^{i+1} x_i \triangleright_{\lambda_i} \gamma_{\lambda_1, \dots, \hat{\lambda}_i, \dots, \lambda_n}(x_1, \dots, \hat{x}_i, \dots, x_{n+1}) \\ &\quad + \sum_{1 \leq i < j \leq n+1} (-1)^{i+j} \gamma_{\lambda_i + \lambda_j, \lambda_1, \dots, \hat{\lambda}_i, \dots, \hat{\lambda}_j, \dots, \lambda_n}([x_i \lambda_i x_j], x_1, \dots, \hat{x}_i, \dots, \hat{x}_j, \dots, x_{n+1}). \end{aligned}$$

It is proved that $\delta^2 = 0$ [BKV99]. Therefore $\{C^*(R, M), \delta\}$ is indeed a cochain complex, whose cohomology is called the cohomology of the Lie conformal algebra R with coefficients in the representation M .

Using the skew-symmetry, we can also write the differential δ as follows:

$$\begin{aligned}
& (\delta\gamma)_{\lambda_1, \dots, \lambda_n}(x_1, \dots, x_{n+1}) \\
&= \sum_{i=1}^n (-1)^{i+1} x_i \triangleright_{\lambda_i} \gamma_{\lambda_1, \dots, \hat{\lambda}_i, \dots, \lambda_n}(x_1, \dots, \hat{x}_i, \dots, x_{n+1}) \\
&\quad + (-1)^{n+1} \gamma_{\lambda_1, \dots, \lambda_{n-1}}(x_1, \dots, x_n) \triangleleft_{\lambda_1 + \dots + \lambda_n} x_{n+1} \\
&\quad + \sum_{1 \leq i < j \leq n+1} (-1)^i \gamma_{\lambda_1, \dots, \hat{\lambda}_i, \dots, \lambda_i + \lambda_j, \dots, \lambda_n}(x_1, \dots, x_{j-1}, [x_i \lambda_i x_j], x_{j+1}, \dots, x_{n+1}).
\end{aligned}$$

This formula was first appeared in [Zha15] as the cohomology of a Leibniz conformal algebras.

Remark 2.8. We use a slightly different notation as in [BKV99; DK98] where they use $\gamma_{\lambda_1, \dots, \hat{\lambda}_i, \dots, \lambda_{n+1}}$ to denote the map γ , but in fact the index λ_{n+1} always doesn't appear in the calculation of the cohomology theory of a Lie conformal algebra, so we use $\gamma_{\lambda_1, \dots, \hat{\lambda}_i, \dots, \lambda_n}$ instead of it.

3. Conformal Lie 2-algebras and 2-term conformal L_∞ -algebras

In this section, we introduced the concept of conformal Lie 2-algebras and 2-term conformal L_∞ -algebras. It is proved that the category of conformal Lie 2-algebras and the category of 2-term conformal L_∞ -algebras are equivalent.

3.1. Conformal Lie 2-algebras

Let's denote the category of $\mathbb{C}[\partial]$ -modules or ∂ -vector spaces by Vect^∂ .

Definition 3.1. A *conformal 2-vector space* is a category in Vect^∂ .

More precisely, a conformal 2-vector space $\mathcal{V}$ is a category with a $\mathbb{C}[\partial]$ -module $\mathcal{V}_0$ and a $\mathbb{C}[\partial]$ -module $\mathcal{V}_1$. The category has source and target maps $s, t: \mathcal{V}_1 \rightarrow \mathcal{V}_0$, an identity-assigning map $i: \mathcal{V}_0 \rightarrow \mathcal{V}_1$, and a composition map $\circ: \mathcal{V}_1 \times_{\mathcal{V}_0} \mathcal{V}_1 \rightarrow \mathcal{V}_1$. These maps are all conformal linear maps. A morphism f from source x to target y is denoted by $f: x \rightarrow y$, where $s(f) = x$ and $t(f) = y$. The notation $i(x)$ is also written as 1_x .

Conformal 2-vector spaces are in one-to-one correspondence with 2-term complexes of $\mathbb{C}[\partial]$ -modules. A 2-term complex of $\mathbb{C}[\partial]$ -modules is a pair of $\mathbb{C}[\partial]$ -modules $\mathcal{C}_1$ and $\mathcal{C}_0$ with a differential between them: $\mathcal{C}_1 \xrightarrow{d} \mathcal{C}_0$. Given a conformal 2-vector space $\mathcal{V}$, $\text{Ker}(s) \xrightarrow{t} \mathcal{V}_0$ is a 2-term complex of $\mathbb{C}[\partial]$ -modules. Conversely, any 2-term complex of $\mathbb{C}[\partial]$ -modules $\mathcal{V}_1 \xrightarrow{d} \mathcal{V}_0$ gives rise to a conformal 2-vector space. In this conformal 2-vector space, the set of objects is $\mathcal{C}_0$ and the set of morphisms is $\mathcal{C}_0 \oplus \mathcal{C}_1$. The source map s is given by $s(x, h) = x$, and the target map t is given by $t(x, h) = x + dh$, where $x \in \mathcal{V}_0$ and $h \in \mathcal{V}_1$. The conformal 2-vector space associated to the 2-term complex

$\mathcal{V}_1 \xrightarrow{d} \mathcal{V}_0$ is denoted by $\mathcal{V}$:

$$\begin{array}{ccc} \mathcal{V}_1 & := & \mathcal{C}_0 \oplus \mathcal{C}_1 \\ \mathcal{V} = & \begin{array}{c} \downarrow s \quad \downarrow t \\ \downarrow \end{array} & \\ \mathcal{V}_0 & := & \mathcal{C}_0. \end{array} \quad (16)$$

Definition 3.2. A *conformal Lie 2-algebra* consists of a conformal 2-vector space $\mathcal{L}$ equipped with

- a skew-symmetric conformal sesquilinear functor, the *bracket*, $[\cdot, \cdot]_\lambda: \mathcal{L} \times \mathcal{L} \rightarrow \mathcal{L}[\lambda]$,
- a conformal sesquilinear natural isomorphism, the *Jacobiator*,

$$J_{x_{\lambda_1} y_{\lambda_2} z}: [x_{\lambda_1} [y_{\lambda_2} z]] \longrightarrow [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z] + [y_{\lambda_2} [x_{\lambda_1} z]],$$

such that the following *Jacobiator identity* is satisfied

$$\begin{aligned} & J_{x_{\lambda_1} y_{\lambda_2} [z_{\lambda_3} t]} \left(1 + [y_{\lambda_2} J_{x_{\lambda_1} z_{\lambda_3} t}] \right) \left(J_{[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z_{\lambda_3} t} + J_{y_{\lambda_2} [x_{\lambda_1} z]_{\lambda_1 + \lambda_3} t} + J_{y_{\lambda_2} z_{\lambda_3} [x_{\lambda_1} t]} \right) \\ &= [x_{\lambda_1} J_{y_{\lambda_2} z_{\lambda_3} t}] \left(J_{x_{\lambda_1} [y_{\lambda_2} z]_{\lambda_2 + \lambda_3} t} + J_{x_{\lambda_1} z_{\lambda_3} [y_{\lambda_2} t]} \right) \left([J_{x_{\lambda_1} y_{\lambda_2} z_{\lambda_1 + \lambda_2 + \lambda_3} t}] + 1 + [z_{\lambda_3} J_{x_{\lambda_1} y_{\lambda_2} t}] \right). \end{aligned} \quad (17)$$

We represent the Jacobiator identity in the following commutative diagram, which illustrates the relationship between two methods of utilizing the Jacobiator to rebracket the expression $[x_{\lambda_1} [y_{\lambda_2} [z_{\lambda_3} t]]]$:

$$\begin{array}{ccccc} & & [x_{\lambda_1} [y_{\lambda_2} [z_{\lambda_3} t]]] & & \\ & \swarrow [x_{\lambda_1} J_{y_{\lambda_2} z_{\lambda_3} t}] & & \searrow J_{x_{\lambda_1} y_{\lambda_2} [z_{\lambda_3} t]} & \\ [x_{\lambda_1} [[y_{\lambda_2} z]_{\lambda_2 + \lambda_3} t]] + [x_{\lambda_1} [z_{\lambda_3} [y_{\lambda_2} t]]] & & & & [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} [z_{\lambda_3} t]] + [y_{\lambda_2} [x_{\lambda_1} [z_{\lambda_3} t]]] \\ \downarrow J_{x_{\lambda_1} [y_{\lambda_2} z]_{\lambda_2 + \lambda_3} t} + J_{x_{\lambda_1} z_{\lambda_3} [y_{\lambda_2} t]} & & & & \downarrow 1 + [y_{\lambda_2} J_{x_{\lambda_1} z_{\lambda_3} t}] \\ [[x_{\lambda_1} [y_{\lambda_2} z]]_{\lambda_1 + \lambda_2 + \lambda_3} t] + [[y_{\lambda_2} z]_{\lambda_2 + \lambda_3} [x_{\lambda_1} t]] & & & & [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} [z_{\lambda_3} t]] + [y_{\lambda_2} [[x_{\lambda_1} z]_{\lambda_1 + \lambda_3} t]] \\ + [[x_{\lambda_1} z]_{\lambda_1 + \lambda_3} [y_{\lambda_2} t]] + [z_{\lambda_3} [x_{\lambda_1} [y_{\lambda_2} t]]] & & & & + [y_{\lambda_2} [z_{\lambda_3} [x_{\lambda_1} t]]] \\ \swarrow [J_{x_{\lambda_1} y_{\lambda_2} z_{\lambda_1 + \lambda_2 + \lambda_3} t}] + 1 + [z_{\lambda_3} J_{x_{\lambda_1} y_{\lambda_2} t}] & & & \nwarrow J_{[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z_{\lambda_3} t} + J_{y_{\lambda_2} [x_{\lambda_1} z]_{\lambda_1 + \lambda_3} t} + J_{y_{\lambda_2} z_{\lambda_3} [x_{\lambda_1} t]} & \\ & P = Q & & & \end{array}$$

where P and Q are given by

$$\begin{aligned} P &= [[[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z]_{\lambda_1 + \lambda_2 + \lambda_3} t] + [[y_{\lambda_2} [x_{\lambda_1} z]]_{\lambda_1 + \lambda_2 + \lambda_3} t] + [[y_{\lambda_2} z]_{\lambda_2 + \lambda_3} [x_{\lambda_1} t]] \\ &\quad + [[x_{\lambda_1} z]_{\lambda_1 + \lambda_3} [y_{\lambda_2} t]] + [z_{\lambda_3} [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} t]] + [z_{\lambda_3} [y_{\lambda_2} [x_{\lambda_1} t]]] = Q. \end{aligned}$$

If we drop the skew-symmetric condition for the bracket, we obtain the concept of conformal Leibniz 2-algebras.

Definition 3.3. A *conformal Leibniz 2-algebra* consists of a conformal 2-vector space $\mathcal{L}$ equipped with a conformal sesquilinear functor $[\cdot, \cdot]_\lambda: \mathcal{L} \times \mathcal{L} \rightarrow \mathcal{L}[\lambda]$ and a conformal sesquilinear natural isomorphism

$$J_{x_{\lambda_1} y_{\lambda_2} z}: [x_{\lambda_1} [y_{\lambda_2} z]] \longrightarrow [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z] + [y_{\lambda_2} [x_{\lambda_1} z]],$$

such that the above (17) is satisfied.

In fact, all the results in the following of this section can be easily generalized to the realm of conformal Leibniz 2-algebra, but we omit the details.

Definition 3.4. Given two conformal Lie 2-algebras $(\mathcal{L}, [\cdot_\lambda \cdot], J)$ and $(\mathcal{L}', [\cdot_\lambda \cdot]', J')$, a conformal Lie 2-algebra morphism $F : \mathcal{L} \rightarrow \mathcal{L}'$ consists of:

- a conformal functor (F^0, F^1) from the underlying conformal 2-vector space of $\mathcal{L}$ to that of $\mathcal{L}'$;
- a natural transformation

$$F_\lambda^2(x, y) : [F^0(x)_\lambda F^0(y)]' \longrightarrow F^0([x_\lambda y])$$

such that the following diagram commutes:

$$\begin{array}{ccc} [F^0(x)_\lambda [F^0(y)_\mu F^0(z)]']' & \xrightarrow{[1, F_\mu^2(y, z)]'} & [F^0(x)_\lambda F^0([y_\mu z])] \\ \downarrow J_{F^0(x)_\lambda F^0(y)_\mu F^0(z)} & & \downarrow F_\lambda^2(x, [y_\mu z]) \\ [[F^0(x)_\lambda F^0(y)]'_{\lambda+\mu} F^0(z)]' + [F^0(y)_\mu [F^0(x)_\lambda F^0(z)]']' & & F^0([x_\lambda [y_\mu z]]) \\ \downarrow [F_\lambda^2(x, y), 1]' + [1, F_\lambda^2(x, z)]' & & \downarrow F^0 J_{x_\lambda y_\mu z} \\ [F^0([x_\lambda y])_{\lambda+\mu} F^0(z)]' + [F^0(y)_\mu F^0([x_\lambda z])] & \xrightarrow{F_{\lambda+\mu}^2([x_\lambda y], z) + F_\mu^2(y, [x_\lambda z])} & F^0([x_\lambda y]_{\lambda+\mu} z) + F^0([y_\mu [x_\lambda z]]). \end{array}$$

The identity morphism $\text{id}_\mathcal{L} : \mathcal{L} \rightarrow \mathcal{L}$ is associated with the identity functor and an identity natural transformation $(\text{id}_\mathcal{L})_2$. The composition of a pair of conformal Lie 2-algebra morphisms, denoted by $G \circ F$, is defined as follows: let $\mathcal{L}$, $\mathcal{L}'$ and $\mathcal{L}''$ be conformal Lie 2-algebras. The functor $((G \circ F)^0, (G \circ F)^1)$ is obtained by composing (G^0, G^1) and (F^0, F^1) , while $(G \circ F)^2$ is defined as the following composition:

$$\begin{array}{ccc} [G^0 \circ F^0(x)_\lambda G^0 \circ F^0(y)]'' & & \\ \downarrow G_\lambda^2(F^0(x), F^0(y)) & \searrow (G \circ F)_\lambda^2(x, y) & \\ & G^0 \circ F^0[x_\lambda y] & \\ & \nearrow G^1(F_\lambda^2(x, y)) & \\ G^0[F^0(x)_\lambda F^0(y)]' & & \end{array}$$

It is easy to see that there is a category LieC2Alg with conformal Lie 2-algebras as objects and conformal Lie 2-algebra morphisms as morphisms.

3.2. 2-term conformal L_∞ -algebras

Now we introduced the concept of 2-term conformal L_∞ -algebras. For general theory of L_∞ -algebras with applications, see [BHR10; Kon03; LS93; RW98; SS85].

Definition 3.5. A 2-term conformal L_∞ -algebra $\mathcal{V} = \mathcal{V}_0 \oplus \mathcal{V}_1$ is a complex consisting of the following data:

- two $\mathbb{C}[\partial]$ -modules $\mathcal{V}_0$ and $\mathcal{V}_1$ together with a conformal linear map $d : \mathcal{V}_1 \rightarrow \mathcal{V}_0$,
- a conformal sesquilinear map $\ell_\lambda^2 : \mathcal{V}_i \times \mathcal{V}_j \rightarrow \mathcal{V}_{i+j}[\lambda]$, where $0 \leq i + j \leq 1$,
- a conformal sesquilinear map $\ell_{\lambda_1, \lambda_2}^3 : \mathcal{V}_0 \times \mathcal{V}_0 \times \mathcal{V}_0 \rightarrow \mathcal{V}_1[\lambda_1, \lambda_2]$.

These maps satisfy the following conditions:

- (a) $l_\lambda^2(x, y) + l_{-\theta-\lambda}^2(y, x) = 0$,
- (b) $l_\lambda^2(x, h) + l_{-\theta-\lambda}^2(h, x) = 0$,
- (c) $l_\lambda^2(h, k) = 0$,
- (d) $l_{\lambda_1, \lambda_2}^3(x, y, z)$ is totally skew symmetric,
- (e) $d(l_\lambda^2(x, h)) = l_\lambda^2(x, dh)$,
- (f) $l_\lambda^2(dh, k) = l_\lambda^2(h, dk)$,
- (g) $d(l_{\lambda_1, \lambda_2}^3(x, y, z)) = l_{\lambda_1}^2(x, l_{\lambda_2}^2(y, z)) - l_{\lambda_1+\lambda_2}^2(l_{\lambda_1}^2(x, y), z) - l_{\lambda_2}^2(y, l_{\lambda_1}^2(x, z))$,
- (h) $l_{\lambda_1, \lambda_2}^3(x, y, dh) = l_{\lambda_1}^2(x, l_{\lambda_2}^2(y, h)) - l_{\lambda_1+\lambda_2}^2(l_{\lambda_1}^2(x, y), h) - l_{\lambda_2}^2(y, l_{\lambda_1}^2(x, h))$,
- (i) $\delta l_{\lambda_1, \lambda_2, \lambda_3}^3(x, y, z, t)$
 $= l_{\lambda_1}^2(x, l_{\lambda_2, \lambda_3}^3(y, z, t)) - l_{\lambda_2}^2(y, l_{\lambda_1, \lambda_3}^3(x, z, t)) + l_{\lambda_3}^2(z, l_{\lambda_1, \lambda_2}^3(x, y, t))$
 $+ l_{\lambda_1+\lambda_2+\lambda_3}^2(l_{\lambda_1, \lambda_2}^3(x, y, z), t) - l_{\lambda_1+\lambda_2, \lambda_3}^2(l_{\lambda_1}^2(x, y), z, t)$
 $- l_{\lambda_2, \lambda_1+\lambda_3}^2(y, l_{\lambda_1}^2(x, z), t) - l_{\lambda_2, \lambda_3}^2(y, z, l_{\lambda_1}^2(x, t)) + l_{\lambda_1, \lambda_2+\lambda_3}^2(x, l_{\lambda_2}^2(y, z), t)$
 $+ l_{\lambda_1, \lambda_3}^2(x, z, l_{\lambda_2}^2(y, t)) - l_{\lambda_1, \lambda_2}^2(x, y, l_{\lambda_3}^2(z, t))$
 $= 0$,

for all $x, y, z, t \in \mathcal{V}_0$ and $h, k \in \mathcal{V}_1$.

Definition 3.6. Let $(\mathcal{V}; d, l_\lambda^2, l_{\lambda_1, \lambda_2}^3)$ and $(\mathcal{V}'; d', l'_\lambda{}^2, l'_{\lambda_1, \lambda_2}{}^3)$ be two 2-term conformal L_∞ -algebras. A CL_∞ -morphism $f = (f^0, f^1, f^2)$ from $\mathcal{V}$ to $\mathcal{V}'$ consists of conformal linear maps $f^0 : \mathcal{V}_0 \rightarrow \mathcal{V}'_0$, $f^1 : \mathcal{V}_1 \rightarrow \mathcal{V}'_1$ and $f^2 : \mathcal{V}_0 \times \mathcal{V}_0 \rightarrow \mathcal{V}'_1[\lambda]$, such that the following equalities hold for all $x, y, z \in \mathcal{V}_0, a \in \mathcal{V}_1$,

- (i) $f^0 d = d' f^1$,
- (ii) $f^0 l_\lambda^2(x, y) - l'_\lambda{}^2(f^0(x), f^0(y)) = d' f_\lambda^2(x, y)$,
- (iii) $f^1 l_\lambda^2(x, a) - l'_\lambda{}^2(f^0(x), f^1(a)) = f_\lambda^2(x, da)$,
- (iv) $f^1(l_{\lambda_1, \lambda_2}^3(x, y, z)) - l'_{\lambda_1, \lambda_2}{}^3(f^0(x), f^0(y), f^0(z))$
 $= f_{\lambda_1}^2(x, l_{\lambda_2}^2(y, z)) - f_{\lambda_1+\lambda_2}^2(l_{\lambda_1}^2(x, y), z) - f_{\lambda_2}^2(y, l_{\lambda_1}^2(x, z))$
 $+ l'_{\lambda_1}{}^2(f^0(x), f_{\lambda_2}^2(y, z)) - l'_{\lambda_1+\lambda_2}{}^2(f_{\lambda_1}^2(x, y), f^0(z)) - l'_{\lambda_2}{}^2(f^0(y), f_{\lambda_1}^2(x, z))$.

If $f^2 = 0$, the CL_∞ -morphisms f is called a strict CL_∞ -morphisms.

Let $f : \mathcal{V} \rightarrow \mathcal{V}'$ and $g : \mathcal{V}' \rightarrow \mathcal{V}''$ be two CL_∞ -morphisms, then their composition $g \circ f : \mathcal{V} \rightarrow \mathcal{V}''$ is a CL_∞ -morphism defined as $(g \circ f)^0 = g^0 \circ f^0$, $(g \circ f)^1 = g^1 \circ f^1$ and

$$(g \circ f)_\lambda^2(x, y) = g_\lambda^2(f^0(x) f^0(y)) + g^1(f_\lambda^2(x, y)).$$

The identity CL_∞ -morphism $1_\mathcal{V} : \mathcal{V} \rightarrow \mathcal{V}$ has the identity chain map together with $(1_\mathcal{V})_2 = 0$.

There is a category $2CL_\infty$ with 2-term conformal L_∞ -algebras as objects and CL_∞ -morphisms as morphisms.

3.3. Equivalence

Now we establish the equivalence between the category of conformal Lie 2-algebras and 2-term conformal L_∞ -algebras.

Theorem 3.7. *The categories $2CL_\infty$ and $LieC2Alg$ are equivalent.*

Proof. First, we show how to construct a conformal Lie 2-algebra from a 2-term conformal L_∞ -algebra.

Let $\mathcal{V} = (\mathcal{V}_1 \xrightarrow{d} \mathcal{V}_0, l_\lambda^2, l_{\lambda_1, \lambda_2}^3)$ be a 2-term conformal L_∞ -algebra, we consider the conformal 2-vector space $\mathcal{L} = (\mathcal{V}_0 \oplus \mathcal{V}_1 \rightrightarrows \mathcal{V}_0)$ given by (16), that is $\mathcal{L}$ has objects $\mathcal{L}_0 = \mathcal{V}_0$, $\mathcal{L}_1 = \mathcal{V}_0 \oplus \mathcal{V}_1$ and morphisms $f: x \rightarrow y$ in $\mathcal{L}_1$ by $f = (x, h)$ where $x \in \mathcal{V}_0$ and $h \in \mathcal{V}_1$. The source, target, and identity-assigning maps in $\mathcal{L}$ are given by

$$\begin{aligned} s(f) &= s(x, h) = x, \\ t(f) &= t(x, h) = x + dh, \\ i(x) &= (x, 0), \end{aligned}$$

and we have $t(f) - s(f) = dh$.

Now we define a skew-symmetric conformal sesquilinear functor on $\mathcal{L}$ by

$$[(x, h)_\lambda(y, k)] = (l_\lambda^2(x, y), l_\lambda^2(x, k) + l_\lambda^2(h, y) + l_\lambda^2(dh, k)),$$

for all $(x, h), (y, k) \in \mathcal{V}_0 \oplus \mathcal{V}_1$. The Jacobiator is given as follows

$$J_{x_{\lambda_1} y_{\lambda_2} z} := ([x_{\lambda_1} [y_{\lambda_2} z]], l_{\lambda_1, \lambda_2}^3(x, y, z)),$$

where $x, y, z \in \mathcal{V}_0$. Then by Condition (g), we have $J_{x_{\lambda_1} y_{\lambda_2} z}$ is a morphism from source $[x_{\lambda_1} [y_{\lambda_2} z]]$ to target $[[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z] + [y_{\lambda_2} [x_{\lambda_1} z]]$.

Next we show that $J_{x_{\lambda_1} y_{\lambda_2} z}$ is a natural isomorphism. We only check naturality in the third variable, the other two cases are similar. Let $f: z \rightarrow z'$. Then, $J_{x_{\lambda_1} y_{\lambda_2} z}$ is natural in z if the following diagram commutes:

$$\begin{array}{ccc} [x_{\lambda_1} [y_{\lambda_2} z]] & \xrightarrow{[[1_x, 1_y], f]} & [x_{\lambda_1} [y_{\lambda_2} z']] \\ \downarrow J_{x_{\lambda_1} y_{\lambda_2} z} & & \downarrow J_{x_{\lambda_1} y_{\lambda_2} z'} \\ [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z] + [y_{\lambda_2} [x_{\lambda_1} z]] & \xrightarrow{[1_x, [1_y, f]] + [[1_x, f], 1_y]} & [[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z'] + [y_{\lambda_2} [x_{\lambda_1} z']] \end{array}$$

Thus we only need to show

$$\begin{aligned} &([x_{\lambda_1} [y_{\lambda_2} z]], l_{\lambda_1, \lambda_2}^3(x, y, z') + [x_{\lambda_1} [y_{\lambda_2} h]]) \\ &= ([x_{\lambda_1} [y_{\lambda_2} z]], [x_{\lambda_1} y]_{\lambda_1 + \lambda_2} h] + [y_{\lambda_2} [x_{\lambda_1} h]] + l_{\lambda_1, \lambda_2}^3(x, y, z)). \end{aligned}$$

This holds by condition (h) in Definition 3.5.

From condition (i) in Definition 3.5, we have

$$\begin{aligned} &l_{\lambda_1}^2(x, l_{\lambda_2, \lambda_3}^3(y, z, t)) - l_{\lambda_2}^2(y, l_{\lambda_1, \lambda_3}^3(x, z, t)) + l_{\lambda_3}^2(z, l_{\lambda_1, \lambda_2}^3(x, y, t)) \\ &+ l_{\lambda_1 + \lambda_2 + \lambda_3}^2(l_{\lambda_1, \lambda_2}^3(x, y, z), t) - l_{\lambda_1 + \lambda_2, \lambda_3}^3(l_{\lambda_1}^2(x, y), z, t) \\ &- l_{\lambda_2, \lambda_1 + \lambda_3}^3(y, l_{\lambda_1}^2(x, z), t) - l_{\lambda_2, \lambda_3}^3(y, z, l_{\lambda_1}^2(x, t)) + l_{\lambda_1, \lambda_2 + \lambda_3}^3(x, l_{\lambda_2}^2(y, z), t) \\ &+ l_{\lambda_1, \lambda_3}^3(x, z, l_{\lambda_2}^2(y, t)) - l_{\lambda_1, \lambda_2}^3(x, y, l_{\lambda_3}^2(z, t)) = 0. \end{aligned}$$

By the skew-symmetry of $l_\lambda^2 = [\cdot, \cdot]_\lambda$, we obtain

$$\begin{aligned} & l_{\lambda_1, \lambda_2}^3(x, y, [z_{\lambda_3} t]) + [y_{\lambda_2} l_{\lambda_1, \lambda_3}^3(x, z, t)] \\ & + l_{\lambda_1 + \lambda_2, \lambda_3}^3([x_{\lambda_1} y], z, t) + l_{\lambda_2, \lambda_1 + \lambda_3}^3(y, [x_{\lambda_1} z], t) + l_{\lambda_2, \lambda_3}^3(y, z, [x_{\lambda_1} t]) \\ & = [x_{\lambda_1} l_{\lambda_2, \lambda_3}^3(y, z, t)] + [z_{\lambda_3} l_{\lambda_1, \lambda_2}^3(x, y, t)] \\ & + [l_{\lambda_1, \lambda_2}^3(x, y, z)_{\lambda_1 + \lambda_2 + \lambda_3} t] + l_{\lambda_1, \lambda_2 + \lambda_3}^3(x, [y_{\lambda_2} z], t) + l_{\lambda_1, \lambda_3}^3(x, z, [y_{\lambda_2} t]). \end{aligned}$$

This is equivalent to Jacobiator identity (17) in Definition 3.2. Thus from a 2-term conformal L_∞ -algebra, we obtain a conformal Lie 2-algebra.

For any CL_∞ -morphism $f = (f^0, f^1, f^2)$ from $\mathcal{V}$ to $\mathcal{V}'$, we construct a conformal Lie 2-algebra morphism $F = T(f)$ from $L = T(\mathcal{V})$ to $L' = T(\mathcal{V}')$ as follows.

Let $F^0 = f^0$, $F^1 = f^0 \oplus f^1$, and F^2 be given by

$$F_\lambda^2(x, y) = ([f^0(x)_\lambda f^0(y)], f_\lambda^2(x, y)).$$

Then $F^2(x, y)$ is a natural isomorphism from $[F^0(x)_\lambda F^0(y)]$ to $F^0[x_\lambda y]$, and $F = (F^0, F^1, F^2)$ is a conformal Lie 2-algebra morphism from $\mathcal{L}$ to $\mathcal{L}'$.

One can also deduce that T preserves the identity CL_∞ -morphisms and the composition of CL_∞ -morphisms. Thus, we obtain a functor from $2CL_\infty$ to $LieC2Alg$.

Conversely, given a conformal Lie 2-algebra $\mathcal{L}$, we define l_λ^2 and $l_{\lambda_1, \lambda_2}^3$ on the 2-term complex $\mathcal{L}_1 \supseteq \ker(s) = \mathcal{V}_1 \xrightarrow{d} \mathcal{V}_0 = \mathcal{L}_0$ by

- $l_1 h = t(h)$ for $h \in \mathcal{V}_1 \subseteq \mathcal{L}_1$,
- $l_\lambda^2(x, y) = [x_\lambda y]$ for $x, y \in \mathcal{V}_0 = \mathcal{L}_0$,
- $l_\lambda^2(x, h) = [1_{x_\lambda} h]$, for $x \in \mathcal{V}_0 = \mathcal{L}_0$ and $h \in \mathcal{V}_1 \subseteq \mathcal{L}_1$,
- $l_\lambda^2(h, k) = 0$ for $h, k \in \mathcal{V}_1 \subseteq \mathcal{L}_1$,
- $l_{\lambda_1, \lambda_2}^3(x, y, z) = p_1 J_{x_{\lambda_1} y_{\lambda_2} z}$ for $x, y, z \in \mathcal{V}_0 = \mathcal{L}_0$, where $p_1 : \mathcal{L}_1 = \mathcal{V}_0 \oplus \mathcal{V}_1 \rightarrow \mathcal{V}_1$ is the projection.

Then one can verify that $(\mathcal{V}_1 \xrightarrow{d} \mathcal{V}_0, l_\lambda^2, l_{\lambda_1, \lambda_2}^3)$ is a 2-term conformal L_∞ -algebra.

Let $F = (F^0, F^1, F^2) : \mathcal{L} \rightarrow \mathcal{L}'$ be a conformal Lie 2-algebra morphism, and $S(\mathcal{L}) = \mathcal{V}$, $S(\mathcal{L}') = \mathcal{V}'$. Define $S(F) = f = (f^0, f^1, f^2)$ as follows. Let $f^0 = F^0$, $f^1 = F^1|_{\mathcal{V}_1 = \text{Ker}(s)}$ and define f^2 by

$$f_\lambda^2(x, y) = F_\lambda^2(x, y) - i(s(F_\lambda^2(x, y))).$$

It is easy to see that f is a $2CL_\infty$ -algebra morphism. Furthermore, S also preserves the identity morphisms and the composition of morphisms. Thus, we obtain a functor from $LieC2Alg$ to CL_∞ .

One show that there are natural isomorphisms $\alpha : T \circ S \rightarrow 1_{LieC2Alg}$ and $\beta : S \circ T \rightarrow 1_{2CL_\infty}$. This complete the proof. $\square$

4. Construction of conformal Lie 2-algebras

In this section, we present concrete examples and special cases of conformal Lie 2-algebras. These include Lie conformal algebras with 3-cocycles, crossed modules of Lie conformal algebras, string Lie conformal algebras, and conformal omni-Lie algebras. Finally, we demonstrate that conformal Lie 2-algebras can be derived from any Leibniz conformal algebra.

4.1. Skeletal conformal Lie 2-algebras

In the case of a 2-term conformal L_∞ -algebra with $d = 0$, it is referred to as *skeletal*. From conditions (a) and (g), it can be deduced that $\mathcal{V}_0$ constitutes a Lie conformal algebra. Furthermore, conditions (b) and (h) indicate that $\mathcal{V}_1$ serves as a representation of $\mathcal{V}_0$ through the action established by $x \triangleright_\lambda h := l_\lambda^2(x, h)$. Moreover, condition (i) can be expressed as a 3-cocycle condition in the Lie conformal algebra cohomology of $\mathcal{V}_0$ with respect to $\mathcal{V}_1$ as its values.

Proposition 4.1. *Skeletal 2-term conformal L_∞ -algebras are in one-to-one correspondence with quadruples $(R, M, \triangleright_\lambda, l_{\lambda_1, \lambda_2}^3)$ where R is a Lie conformal algebra, M is a $\mathbb{C}[\partial]$ -module, $\triangleright_\lambda$ is an action of R on M and $l_{\lambda_1, \lambda_2}^3$ is a 3-cocycle on R with coefficients in M .*

Let M be a $\mathbb{C}[\partial]$ -module. A conformal bilinear form on M is a $\mathbb{C}$ -bilinear map $\langle \cdot, \cdot \rangle_\lambda : M \times M \rightarrow \mathbb{C}[\lambda]$ such that

$$\langle \partial v, w \rangle_\lambda = -\lambda \langle v, w \rangle_\lambda = -\langle v, \partial w \rangle_\lambda, \quad \text{for all } v, w \in M.$$

The conformal bilinear form is symmetric if $\langle v, w \rangle_\lambda = \langle w, v \rangle_{-\lambda}$ for all $v, w \in V$. The conformal bilinear form in a Lie conformal algebra R is called invariant if

$$\langle [x_\mu y], z \rangle_\lambda = \langle x, [y_{\lambda-\partial} z] \rangle_\mu = -\langle x, [z_{-\lambda} y] \rangle_\mu,$$

for all $x, y, z \in R$.

Example 4.2. Given a Lie conformal algebra $(R, [\cdot_\lambda \cdot])$ with symmetric invariant conformal bilinear form $\langle \cdot, \cdot \rangle_\lambda$, we construct a 2-term conformal L_∞ -algebra as follows. Let $\mathcal{V}_1 = \mathbb{C}$, $\mathcal{V}_0 = R$, $d = 0$, and define $l_\lambda^2, l_{\lambda_1, \lambda_2}^3$ by

$$l_\lambda^2(x, y) = [x_\lambda y], \quad l_\lambda^2(x, h) = 0, \quad l_{\lambda_1, \lambda_2}^3(x, y, z) = \langle [x_{\lambda_1} y], z \rangle_{\lambda_2}, \quad (18)$$

where $x, y, z \in R, h \in \mathbb{C}$. All the conditions in the Definition 3.5 are satisfied. Thus we obtain a 2-term conformal L_∞ -algebra $(\mathbb{C} \xrightarrow{0} R, l_\lambda^2, l_{\lambda_1, \lambda_2}^3)$. We call it the *string conformal Lie 2-algebra*.

4.2. Strict conformal Lie 2-algebras

Another kind of 2-term conformal L_∞ -algebra is called *strict* if $l_{\lambda_1, \lambda_2}^3 = 0$. These conformal Lie 2-algebras can be characterized using crossed modules of Lie conformal algebras.

Definition 4.3. Let $(R, [\cdot, \cdot]_R)$ and $(H, [\cdot, \cdot]_H)$ be two Lie conformal algebras. A crossed module of Lie conformal algebras is a homomorphism of Lie conformal algebras $\varphi : H \rightarrow R$ together with an action of R on H , denoted by $x \triangleright_\lambda h$, such that

$$\varphi(x \triangleright_\lambda h) = [x_\lambda \varphi(h)]_R, \quad \varphi(h) \triangleright_\lambda k = [h_\lambda k]_H,$$

for all $h, k \in H, x \in R$.

Proposition 4.4. *There is an one-to-one correspondence between strict 2-term conformal L_∞ -algebras and crossed modules of Lie conformal algebras.*

Proof. Let $\mathcal{V}_1 \xrightarrow{d} \mathcal{V}_0$ be a 2-term conformal L_∞ -algebra with $l_{\lambda_1, \lambda_2}^3 = 0$. We construct Lie conformal algebras on $R = \mathcal{V}_0$ and $H = \mathcal{V}_1$ as follows. The bracket on R and H are defined by

$$\begin{aligned} [h_\lambda k]_H &:= l_\lambda^2(dh, k), \quad \forall h, k \in H = \mathcal{V}_1, \\ [x_\lambda y]_R &:= l_\lambda^2(x, y), \quad \forall x, y \in R = \mathcal{V}_0. \end{aligned}$$

By condition (a) and (g) in Definition 3.5, it is easy to see that $[\cdot_\lambda \cdot]_R$ satisfies the Jacobi identity. By (h), we have

$$\begin{aligned} [h_\lambda [k_\mu l]] - [[h_\lambda k]_{\lambda+\mu} l] - [k_\mu [h, {}_\lambda l]] \\ = l_\lambda^2(dh, l_\mu^2(dk, l)) - l_{\lambda+\mu}^2(dl_\lambda^2(dh, k), l) - l_\mu^2(dk, l_\lambda^2(dh, l)) \\ = l_\lambda^2(dh, l_\mu^2(dk, l)) - l_{\lambda+\mu}^2(l_\lambda^2(dh, dk), l) - l_\mu^2(dk, l_\lambda^2(dh, l)) \\ = 0. \end{aligned}$$

Thus $[\cdot_\lambda \cdot]_H$ satisfies the Jacobi identity. Now let $\varphi = d$, then by (e), we have

$$\varphi([h_\lambda k]_H) = d(l_\lambda^2(dh, k)) = l_\lambda^2(dh, dk) = [\varphi(h)_\lambda \varphi(k)]_R,$$

which implies that φ is a homomorphism of Lie conformal algebras.

Now define the maps of $\triangleright_\lambda : R \times H \rightarrow H[\lambda]$ by

$$x \triangleright_\lambda h := l_\lambda^2(x, h), \quad \forall x \in R, h \in H,$$

which is an action of R on H and it is easy to check that

$$\begin{aligned} \varphi(x \triangleright_\lambda h) &= d(l_\lambda^2(x, h)) = l_\lambda^2(x, dh) = [x_\lambda \varphi(h)]_R, \\ \varphi(h) \triangleright_\lambda k &= l_\lambda^2(dh, k) = [h_\lambda k]_H. \end{aligned}$$

Therefore, we obtain a crossed module of Lie conformal algebras.

Conversely, a crossed module of Lie conformal algebras gives rise to a 2-term conformal L_∞ -algebra with $d = \varphi$, $\mathcal{V}_0 = R$ and $\mathcal{V}_1 = H$, where the brackets are given by

$$\begin{aligned} l_\lambda^2(x, y) &:= [x_\lambda y]_R, \quad \forall x, y \in R, \\ l_\lambda^2(x, h) &:= x \triangleright_\lambda h, \quad \forall h \in H, \\ l_\lambda^2(h, k) &:= 0. \end{aligned}$$

The crossed module conditions give various conditions for 2-term conformal L_∞ -algebras with $l_{\lambda_1, \lambda_2}^3 = 0$. $\square$

Given a crossed module (R, H, φ) , we consider $\mathfrak{g} = \ker(\varphi)$ and $M = \ker(\varphi)$. The Lie conformal algebra structure on H induces a Lie conformal algebra structure on $\mathfrak{g}$. Moreover, the action of H on R induces an action of $\mathfrak{g}$ on M . Hence a crossed module yields an exact sequence

$$0 \longrightarrow M \xrightarrow{i} R \xrightarrow{\varphi} H \xrightarrow{\pi} \mathfrak{g} \longrightarrow 0,$$

where $\pi : H \rightarrow \mathfrak{g}$ is the projection map. Two crossed modules (R, H, φ) and (R', H', φ') over $\mathfrak{g}$ with kernel M are said to be equivalent if there is a morphism of crossed modules $(R, H, \varphi) \rightarrow (R', H', \varphi')$ which induces identity maps on $\mathfrak{g}$ and M . Let $\text{Cross}(\mathfrak{g}, M)$ be the equivalence classes of such crossed modules.

Proposition 4.5. *The equivalence classes of crossed modules are in one-to-one correspondence with the elements of third cohomology group of $\mathfrak{g}$ with coefficients in M .*

Proof. Given a crossed module extension

$$0 \longrightarrow M \xrightarrow{i} R \xrightarrow{\varphi} H \xrightarrow{\pi} \mathfrak{g} \longrightarrow 0.$$

Choose $\mathbb{C}$ -linear sections $s : \mathfrak{g} \rightarrow H$ with $\pi s = \text{id}$ and $q : \text{Im}(\partial) \rightarrow R$ with $\partial q = \text{id}$. Since π is a Lie conformal algebra homomorphism, for all $x, y \in \mathfrak{g}$, we have

$$\pi([s(x)_\lambda s(y)] - s[x_\lambda y]) = [\pi s(x)_\lambda \pi s(y)] - \pi s[x_\lambda y] = [x_\lambda y] - [x_\lambda y] = 0.$$

This shows that $[s(x)_\lambda s(y)] - s[x_\lambda y] \in \ker(\pi) = \text{Im}(\partial)$. Take

$$g_\lambda(x, y) = q([s(x)_\lambda s(y)] - s[x_\lambda y]) \in R.$$

Define a map $\theta_{\mathcal{E}, s, q} : \mathfrak{g}^{\otimes 3} \rightarrow M[\lambda]$ by

$$\begin{aligned} \theta_{\mathcal{E}, s, q}(x, y, z) &= [s(x)_\lambda g_\mu(y, z)] - [g_\lambda(x, y)_{\lambda+\mu} s(z)] - [s(y)_\nu g_\lambda(x, z)] \\ &\quad - g_{\lambda+\mu}([x_\lambda y], z) + g_\lambda(x, [y_\mu z]) + g_{-\partial-\mu}([x_\lambda z], y). \end{aligned}$$

Since ∂ is a map of H -modules, it follows that $\partial(\theta_{\mathcal{E}, s, q}(x, y, z)) = 0$. Thus we obtain $\theta_{\mathcal{E}, s, q}(x, y, z) \in \ker(\partial) = M$. Therefore $\theta_{\mathcal{E}, s, q} : \mathfrak{g}^{\otimes 3} \rightarrow M[\lambda]$ is well defined. Moreover, the map $\theta_{\mathcal{E}, s, q}$ is a 3-cocycle ($\delta(\theta_{\mathcal{E}, s, q}) = 0$) in the cohomology of $\mathfrak{g}$ with coefficients in M . It follows that equivalence classes of crossed modules gives the same cohomology 3-cocycle classes. The proof is similar as the Lie algebra case, see Wagemann [Wag06] for more details. $\square$

Let R be a Lie conformal algebra, recall that a map $D \in \text{Cend}(R)$ is called a conformal derivation [GT14] if

$$D_\lambda[x_\mu y] = [(D_\lambda x)_{\lambda+\mu} y] + [x_\mu (D_\lambda y)].$$

For example, the adjoint action $(\text{ad } x)_\lambda$ defined by

$$(\text{ad } x)_\lambda(y) = [x_\lambda y]$$

from R to itself is a conformal derivation. This is called the inner conformal derivation.

Denote by $\text{Der}(R)$ the vector space spanned by all conformal derivations. We find that $\text{Der}(R)$ becomes a Lie conformal algebra under the bracket

$$[D_\lambda D'] = D_\lambda D' - D'_{-\partial-\lambda} D,$$

equivalently,

$$[D_\lambda D']_{\lambda'} = D_\lambda D'_{\lambda'-\lambda} - D'_{\lambda'-\lambda} D_\lambda,$$

where D, D' are conformal derivations of R .

For an operation ω on R such that $\omega_\lambda : R \times R \rightarrow R[\lambda]$, we define the adjoint operator $\text{ad}_\omega : R \rightarrow \text{Cend}(R)$ by

$$\text{ad}_{\omega_\lambda}(x)(y) = \omega_\lambda(x, y) \in R, \quad \forall x, y \in R.$$

Then the graph of the adjoint operator

$$\mathcal{F}_\omega := \{\text{ad}_\omega x + x \mid \forall x \in R\} \subset \mathcal{E} := \text{Cend}(R) \oplus R$$

is a subspace of $\mathcal{E}$.

Proposition 4.6. *D is a conformal derivation of R if and only if $\mathcal{F}_\omega$ is an invariant subspace of D , that is $D \circ \mathcal{F}_\omega \subseteq \mathcal{F}_\omega$, where the action D on $\mathcal{F}_\omega$ is defined by*

$$D_\lambda \circ (\text{ad}_{\omega_\mu}(x) + x) = [D_\lambda \text{ad}_{\omega_\mu}(x)] + D_\lambda x, \quad (19)$$

for all $\text{ad}_{\omega_\mu}(x) + x \in \mathcal{F}_{\omega_\mu}$.

Proof. The right hand side of (19) is contained in $\mathcal{F}_{\omega_\mu}$ if and only if

$$[D_\lambda \text{ad}_{\omega_\mu}(x)] = \text{ad}_{\omega_\mu}(D_\lambda x),$$

this is equivalent to

$$\begin{aligned} D_\lambda \text{ad}_{\omega_\mu}(x)(y) - \text{ad}_{\omega_\mu}(x)D(y) - \text{ad}_{\omega_\mu}(D_\lambda x)(y) &= D_\lambda[x_\mu y] - [x_\mu(D_\lambda y)] - [(D_\lambda x)_{\lambda+\mu}y] \\ &= 0. \end{aligned}$$

Thus D is a conformal derivation if and only if $D \circ \mathcal{F}_\omega \subseteq \mathcal{F}_\omega$. $\square$

We call the set of $D \in \text{Cend}(R)$ such that $D \circ \mathcal{F}_\omega \subseteq \mathcal{F}_\omega$ the normalizer of $\mathcal{F}_\omega$, which is denoted by $N(\mathcal{F}_\omega)$.

Proposition 4.7. *Let D and D' be two conformal derivations. Then $[D_\lambda D']$ is also a conformal derivation. Thus we have $\text{Der}(R) = N(\mathcal{F}_\omega)$ is a conformal Lie subalgebra of $\text{Cend}(R)$.*

Proof. Let $D \circ \mathcal{F}_\omega \subseteq \mathcal{F}_\omega$ and $D' \circ \mathcal{F}_\omega \subseteq \mathcal{F}_\omega$, then

$$[D_\lambda \text{ad}_{\omega_\mu}(x)] = \text{ad}_{\omega_\mu}(D_\lambda x), \quad [D'_{\lambda'} \text{ad}_{\omega_\mu}(x)] = \text{ad}_{\omega_\mu}(D'_{\lambda'} x).$$

By the Jacobi identity of $\text{Cend}(R)$, we have

$$\begin{aligned} [[D_\lambda D']_{\lambda+\lambda'} \text{ad}_{\omega_\mu}(x)] &= [D_\lambda [D'_{\lambda'} \text{ad}_{\omega_\mu}(x)]] - [D'_{\lambda'} [D_\lambda \text{ad}_{\omega_\mu}(x)]] \\ &= [D_\lambda \text{ad}_{\omega_\mu}(D'_{\lambda'} x)] - [D'_{\lambda'} \text{ad}_{\omega_\mu}(D_\lambda x)] \\ &= \text{ad}_{\omega_\mu}(D_\lambda D'_{\lambda'} x) - \text{ad}_{\omega_\mu}(D'_{\lambda'} D_\lambda x) \\ &= \text{ad}_{\omega_\mu}([D_\lambda D']_{\lambda+\lambda'} x). \end{aligned}$$

This is equivalent to $[D_\lambda D'] \circ \mathcal{F}_\omega \subseteq \mathcal{F}_\omega$, so the bracket in $\text{Der}(R)$ is closed. Thus $\text{Der}(R)$ is a Lie conformal algebra as a conformal Lie subalgebra of $\text{Cend}(R)$. $\square$

Proposition 4.8. *Let R be a Lie conformal algebra, $\text{Der}(R)$ and $\text{Inn}(R)$ be the set of their conformal derivations and inner conformal derivations. Then we obtain a crossed module $i : \text{Inn}(R) \rightarrow \text{Der}(R)$, with $\text{Der}(R)$ acting $\text{Inn}(R)$ by $D \triangleright_\lambda \text{ad}_x = \text{ad}_{D_\lambda x}$.*

4.3. Conformal omni-Lie algebras

The notion of omni-Lie algebras was generalized to omni-Lie superalgebras in [ZL14]. In this subsection, we introduce the concept of conformal omni-Lie algebras and construct 2-term conformal L_∞ -algebras from them.

Let R be a Lie conformal algebra and M be a left R -module. We define an operation $\circ$ on $\mathcal{E} := R \oplus M$ by

$$(x + u) \circ_\lambda (y + v) = [x_\lambda y] + x \triangleright_\lambda v, \quad (20)$$

for all $x, y \in R$ and $u, v \in M$. Then we have

Proposition 4.9. $(\mathcal{E}, \circ_\lambda)$ is a Leibniz conformal algebra.

Proof. We have to check the Leibniz identity for the operation $\circ_\lambda$. Let $e_1 = x + u, e_2 = y + v, e_3 = z + w$. Then we have

$$\begin{aligned} & \{e_1 \circ_\lambda e_2\} \circ_{\lambda+\mu} e_3 - e_1 \circ_\lambda \{e_2 \circ_\mu e_3\} - e_2 \circ_\mu \{e_1 \circ_\lambda e_3\} \\ &= ([x_\lambda y] + x \triangleright_\lambda v) \circ_{\lambda+\mu} (z + w) - (x + u) \circ_\lambda ([y_\mu z] + y \triangleright_\mu w) \\ &\quad - (y + v) \circ_\mu ([x \triangleright_\lambda w] + z \triangleright_\lambda w) \\ &= [[x_\lambda y]_{\lambda+\mu} z] - [x_\lambda [y_\mu z]] - [y_\mu [x_\lambda z]] \\ &\quad + [x_\lambda y] \triangleright_{\lambda+\mu} w - x \triangleright_\lambda (y \triangleright_\mu w) - y \triangleright_\mu (x \triangleright_\lambda w) \\ &= 0. \end{aligned}$$

The right hand side of the above equation is zero because R is a Lie conformal algebra acting on M . $\square$

We call this type of Leibniz conformal algebra structure on $\mathcal{E}$ the *hemisemidirect product* of R with M as in [KW01], denote it by $R \ltimes_H M$. Note that the operation $\circ_\lambda$ is not skew-symmetry. The skew-symmetrized bracket of it is:

$$[(x + u)_\lambda (y + v)] := [x_\lambda y] + \frac{1}{2} (x \triangleright_\lambda v - y \triangleright_{-\partial-\lambda} u). \quad (21)$$

This is called *demisemidirect product* of R with M , denoted by $R \ltimes_D M$.

For $e_1 = x + u, e_2 = y + v, e_3 = z + w \in \mathcal{E}$, denote by

$$\tilde{J}(e_1, e_2, e_3) := [[e_{1\lambda_1} [e_{2\lambda_2} e_3]]] - [[e_{2\lambda_2} [e_{1\lambda_1} e_3]]] - [[[e_{1\lambda_1} e_2]]_{\lambda_1+\lambda_2} e_3].$$

Then $\tilde{J}(e_1, e_2, e_3) \neq 0$. For example, let $e_1 = x, e_2 = y, e_3 = w$, then

$$\tilde{J}(x, y, z) = -\frac{1}{4} [x_{\lambda_1} y] \triangleright_{\lambda_1+\lambda_2} w.$$

Thus $R \ltimes_D M$ can not to be a Lie conformal algebra but a Lie conformal 2-algebra, see Theorem 4.10 as follows.

Now, for a Lie conformal algebra R and a left R -module M , let

$$\mathcal{V}_0 = R \ltimes_D M, \quad \mathcal{V}_1 = M, \quad d = i : M \hookrightarrow R \ltimes_D M,$$

where i is the inclusion map and define

$$l_\lambda^2 = [[\cdot \cdot \cdot]], \quad l_{\lambda_1, \lambda_2}^3 = \tilde{J}.$$

Theorem 4.10. Let R be a Lie conformal algebra and M a left R -module. Then we obtain a nontrivial 2-term conformal L_∞ -algebra $(M \xrightarrow{i} R \ltimes_D M, l_\lambda^2, l_{\lambda_1, \lambda_2}^3)$.

Proof. It can be checked that various conditions in Definition 3.5 hold. First we have

$$\begin{aligned} & [(x + u)_\lambda (y + v)] + [(y + v)_{-\partial-\lambda} (x + u)] \\ &= [x_\lambda y] + \frac{1}{2} (x \triangleright_\lambda v - y_{-\partial-\lambda} u) + [y_{-\partial-\lambda} x] + \frac{1}{2} (y_{-\partial-\lambda} u - x \triangleright_\lambda v) \\ &= [x_\lambda y] + [y_{-\partial-\lambda} x] + \frac{1}{2} (x \triangleright_\lambda v - y_{-\partial-\lambda} u) + \frac{1}{2} (y_{-\partial-\lambda} u - x \triangleright_\lambda v) \\ &= 0. \end{aligned}$$

Thus condition (a) holds.

Let $e_1 = x, e_2 = y, e_3 = z$ where $x, y, z \in R$ and $e_4 = t \in M$, then we have

$$\begin{aligned}
& \llbracket x_{\lambda_1} l_{\lambda_2, \lambda_3}^3(y, z, t) \rrbracket - \llbracket y_{\lambda_2} l_{\lambda_1, \lambda_3}^3(x, z, t) \rrbracket \\
& + \llbracket z_{\lambda_3} l_{\lambda_1, \lambda_2}^3(x, y, t) \rrbracket + \llbracket l_{\lambda_1, \lambda_2}^3(x, y, z)_{\lambda_1 + \lambda_2 + \lambda_3} t \rrbracket - l_{\lambda_1 + \lambda_2, \lambda_3}^3(\llbracket x_{\lambda_1} y \rrbracket, z, t) \\
& - l_{\lambda_2, \lambda_1 + \lambda_3}^3(y, \llbracket x_{\lambda_1} z \rrbracket, t) - l_{\lambda_2, \lambda_3}^3(y, z, \llbracket x_{\lambda_1} t \rrbracket) \\
& + l_{\lambda_1, \lambda_2 + \lambda_3}^3(x, \llbracket y_{\lambda_2} z \rrbracket, t) + l_{\lambda_1, \lambda_3}^3(x, z, \llbracket y_{\lambda_2} t \rrbracket) - l_{\lambda_1, \lambda_2}^3(x, y, \llbracket z_{\lambda_3} t \rrbracket) \\
& = -\frac{1}{8}x \triangleright_{\lambda_1} ([y_{\lambda_2} z]_{\lambda_2 + \lambda_3} \triangleright t) + \frac{1}{8}y \triangleright_{\lambda_2} ([x_{\lambda_1} z]_{\lambda_1 + \lambda_3} \triangleright t) \\
& - \frac{1}{8}z \triangleright_{\lambda_3} ([x, y]_{\lambda_1 + \lambda_2} \triangleright t) + 0 + \frac{1}{4}[\llbracket x_{\lambda_1} y \rrbracket_{\lambda_1 + \lambda_2} z]_{\lambda_1 + \lambda_2 + \lambda_3} t \\
& - \frac{1}{4}[\llbracket x_{\lambda_1} z \rrbracket_{\lambda_1 + \lambda_3} y]_{\lambda_1 + \lambda_2 + \lambda_3} t + \frac{1}{4}[\llbracket y_{\lambda_2} z \rrbracket_{\lambda_2 + \lambda_3} x]_{\lambda_1 + \lambda_2 + \lambda_3} t \\
& + \frac{1}{8}[y_{\lambda_2} z]_{\lambda_2 + \lambda_3} (x \triangleright_{\lambda_1} t) - \frac{1}{8}[x_{\lambda_1} z]_{\lambda_1 + \lambda_3} (y \triangleright_{\lambda_2} t) + \frac{1}{8}[x_{\lambda_1} y]_{\lambda_1 + \lambda_2} (z \triangleright_{\lambda_3} t) \\
& = -\frac{3}{8}(\llbracket x_{\lambda_1} [y_{\lambda_2} z] \rrbracket - \llbracket [x_{\lambda_1} y]_{\lambda_1 + \lambda_2} z \rrbracket - [y_{\lambda_2} [x_{\lambda_1} z]]_{\lambda_1 + \lambda_2 + \lambda_3} t) \\
& = 0.
\end{aligned}$$

Thus condition (i) holds. The other conditions can be checked similarly. $\square$

Furthermore, we define a symmetric bilinear form on $\mathcal{E}$ with values in $M[\lambda]$ by

$$\langle x + u, y + v \rangle_{\lambda} := \frac{1}{2}(x \triangleright_{\lambda} v + y \triangleright_{\lambda} u). \quad (22)$$

The triple $(\mathcal{E}, \llbracket \cdot_{\lambda} \cdot \rrbracket, \langle \cdot, \cdot \rangle_{\lambda})$ is called a *conformal omni-Lie algebra*.

When $R = \text{Cend}(M)$, all possible Lie conformal algebra structures on M can be characterized by means of the conformal omni-Lie algebras.

For an operation $\omega_{\lambda} : M \times M \rightarrow M[\lambda]$, we define the adjoint operator

$$\text{ad}_{\omega_{\lambda}} : M \text{Cend}(M), \quad \text{ad}_{\omega_{\lambda}}(x)(y) = \omega_{\lambda}(x, y) \in M[\lambda],$$

where $x, y \in M$. Then the graph of the adjoint operator:

$$\mathcal{F}_{\omega} = \{\text{ad}_{\omega} x + x \mid \forall x \in M\} \subset \mathcal{E} = \text{Cend}(M) \oplus M$$

is a subspace of $\mathcal{E}$. Denote $\mathcal{F}_{\omega}^{\perp}$ the orthogonal complement of $\mathcal{F}_{\omega}$ in $\mathcal{E}$ with respect to the symmetric bilinear form $\langle \cdot, \cdot \rangle$ on $\mathcal{E}$ given in (22).

Proposition 4.11. *With the above notations, (M, ω_{λ}) is a Lie conformal algebra if and only if its graph $\mathcal{F}_{\omega}$ is maximal isotropic, i.e. $\mathcal{F}_{\omega} = \mathcal{F}_{\omega}^{\perp}$, and is closed with respect to the bracket $\llbracket \cdot_{\lambda} \cdot \rrbracket$.*

Proof. If ω_{λ} is skew symmetric, i.e. $\omega_{\lambda}(x, y) + \omega_{-\partial - \lambda}(y, x) = 0$, then

$$\begin{aligned}
\langle \text{ad}_{\omega_{\lambda}}(x) + x, \text{ad}_{\omega_{-\partial - \lambda}}(y) + y \rangle_{\lambda} &= \frac{1}{2}(\text{ad}_{\omega_{\lambda}}(x)y + \text{ad}_{\omega_{-\partial - \lambda}}(y)x) \\
&= \frac{1}{2}(\omega_{\lambda}(x, y) + \omega_{-\partial - \lambda}(y, x)).
\end{aligned}$$

This means that ω_{λ} is skew-symmetric if and only if its graph is isotropic, i.e. $\mathcal{F}_{\omega} \subseteq \mathcal{F}_{\omega}^{\perp}$. Moreover, by dimension analysis, we have $\mathcal{F}_{\omega}$ is maximal isotropic.

Next let $[x_\lambda y] := \omega_\lambda(x, y)$, we shall check that the Jacobi identity on M is satisfied if and only if $\mathcal{F}_\omega$ is closed under bracket (21) on $\mathcal{E}$. In fact,

$$\begin{aligned} [[\text{ad}_{\omega_\lambda}(x) + x, \text{ad}_{\omega_\mu}(y) + y]] &= [\text{ad}_{\omega_\lambda}(x), \text{ad}_{\omega_\mu}(y)] + \frac{1}{2}(\text{ad}_{\omega_\lambda}(x)y - \text{ad}_{\omega_{-\partial-\mu}}(y)x) \\ &= [\text{ad}_{\omega_\lambda}(x), \text{ad}_{\omega_\mu}(y)] + \frac{1}{2}(\omega_\lambda(x, y) - \omega_{-\partial-\lambda}(y, x)) \\ &= [\text{ad}_{\omega_\lambda}(x), \text{ad}_{\omega_\mu}(y)] + \omega_\lambda(x, y). \end{aligned}$$

Thus this bracket is closed if and only if

$$[\text{ad}_{\omega_\lambda}(x), \text{ad}_{\omega_\mu}(y)] = \text{ad}_{\omega_{\lambda+\mu}}(\omega_\lambda(x, y)).$$

In this case, for all $z \in M$, we have

$$\begin{aligned} &[\text{ad}_{\omega_\lambda}(x), \text{ad}_{\omega_\mu}(y)](z) - \text{ad}_{\omega_{\lambda+\mu}}(\omega_\lambda(x, y))(z) \\ &= \text{ad}_{\omega_\lambda}(x) \text{ad}_{\omega_\mu}(y)(z) - \text{ad}_{\omega_\mu}(y) \text{ad}_{\omega_\lambda}(x)(z) - \text{ad}_{\omega_{\lambda+\mu}}(\omega_\lambda(x, y))(z) \\ &= \text{ad}_{\omega_\lambda}(x) \omega_\mu(y, z) - \text{ad}_{\omega_\mu}(y) \omega_\lambda(x, z) - \omega_{\lambda+\mu}(\omega_\lambda(x, y), z) \\ &= \omega_\lambda(x, \omega_\mu(y, z)) - \omega_\mu(y, \omega_\lambda(x, z)) - \omega_{\lambda+\mu}(\omega_\lambda(x, y), z) \\ &= [x_\lambda[y_\mu z]] - [y_\mu[x_\lambda z]] - [[x_\lambda y]_{\lambda+\mu} z] \\ &= 0. \end{aligned}$$

This is exactly the Jacobi identity on M . □

We define a Dirac structure of $\text{Cend}(M) \oplus M$ to be any maximal isotropic subspace $L \subseteq \text{Cend}(M) \oplus M$ which is closed under the bracket operation. Then we obtain that (M, ω_λ) is a Lie conformal algebra if and only if $\mathcal{F}_{\omega_\lambda}$ is a Dirac structure of the conformal omni-Lie algebra $R \oplus M$.

For a general characterization for all Dirac structures of $\mathcal{E}$, we adapt the theory of characteristic pairs developed in [Liu00]. One can prove the following result.

Proposition 4.12. *There is a one-to-one correspondence between Dirac structures of the conformal omni-Lie algebra $(\text{Cend}(M) \oplus M, \llbracket \cdot, \cdot \rrbracket_\lambda, \langle \cdot, \cdot \rangle_\lambda)$ and Lie conformal algebra structures on subspaces of M .*

4.4. Skew-symmetrization of Leibniz conformal algebras

From above subsection, we have seen that a conformal omni-Lie algebra $\mathcal{E} = R \ltimes_D M$ is in fact the skew-symmetrization of Leibniz conformal algebra $R \ltimes_H M$ and M is the left center of $\mathcal{E}$. This observation prompts the following question: Can we derive a conformal Lie 2-algebra from any given Leibniz conformal algebra? In the following subsection, we provide a positive answer to this inquiry.

Let L be a Leibniz conformal algebra. We define the *Leibniz kernel* $\text{Ker}(L)$ of L to be the set of elements spanned by

$$\{x \circ_\lambda y + y \circ_{-\partial-\lambda} x \mid \forall x, y \in L\}.$$

Note that if L is a Lie conformal algebra, then $\text{Ker}(L) = 0$, otherwise if L is a non-Lie Leibniz conformal algebra, then $\text{Ker}(L) \neq 0$.

The *left center* $Z^l(L)$ of L is defined by

$$Z^l(L) = \{t \in L \mid t \circ_\lambda x = 0, \forall x \in L\}.$$

It is easy to see that $\text{Ker}(L)$ and $Z^l(L)$ are ideals of L and the quotient algebras $L/\text{Ker}(L)$ and $L/Z^l(L)$ are Lie conformal algebras.

Proposition 4.13. *The Leibniz kernel $\text{Ker}(L)$ is contained in the left center $Z^l(L)$. Thus for any non-Lie Leibniz conformal algebra, the set $Z^l(L)$ is not empty.*

Proof. Let $x \circ_\lambda y + y \circ_{-\partial-\lambda} x \in \text{Ker}(L)$. Then we get

$$\begin{aligned} (x \circ_\lambda y + y \circ_{-\partial-\lambda} x) \circ_{\lambda+\mu} z \\ = x \circ_\lambda (y \circ_\mu z) - y \circ_\mu (x \circ_\lambda z) + y \circ_\mu (x \circ_\lambda z) - x \circ_\lambda (y \circ_\mu z) \\ = 0, \end{aligned}$$

for all $z \in L$. Thus $x \circ_\lambda y + y \circ_{-\partial-\lambda} x \in Z^l(L)$. Therefore $\text{Ker}(L)$ is contained in $Z^l(L)$. $\square$

For a Leibniz conformal algebra $(L, \circ_\lambda)$, since the operation $\circ_\lambda$ is not skew-symmetry, we introduce the following skew-symmetric bracket on L by

$$[x_\lambda y] = \frac{1}{2} (x \circ_\lambda y - y \circ_{-\partial-\lambda} x), \quad \forall x, y \in L, \quad (23)$$

and let $\tilde{J}$ be given by

$$\tilde{J}(x, y, z) := [[x_{\lambda_1} y]_{\lambda_1+\lambda_2} z] - [x_{\lambda_1} [y_{\lambda_2} z]] - [y_{\lambda_2} [x_{\lambda_1} z]]. \quad (24)$$

Now it is easy to prove that

Proposition 4.14. *If $t \in Z^l(L)$, i.e. $t \circ x = 0$ for all $x \in L$, then we have*

$$[x_\lambda t] = \frac{1}{2} x \circ_\lambda t, \quad \tilde{J}(x, y, t) = -\frac{1}{4} (x \circ_{\lambda_1} y) \circ_{\lambda_1+\lambda_2} t.$$

At last, for any non-Lie Leibniz conformal algebra L , we construct nontrivial conformal Lie 2-algebras as follows. Let

$$\mathcal{V}_0 = L, \quad \mathcal{V}_1 = Z^l(L), \quad d = i : Z^l(L) \hookrightarrow L, \quad l_\lambda^2 = [\cdot_\lambda \cdot], \quad l_{\lambda_1, \lambda_2}^3 = \tilde{J}.$$

Theorem 4.15. *With the above notations, from a Leibniz conformal algebra $(L, \circ_\lambda)$, we derive a nontrivial 2-term conformal L_∞ -algebra $(Z^l(L) \xrightarrow{d} L, l_\lambda^2, l_{\lambda_1, \lambda_2}^3)$.*

Proof. By definition of d , l_λ^2 and $l_{\lambda_1, \lambda_2}^3$, it is easy to see that conditions (a)–(h) hold. For condition (i), we verify the case of $x, y, z \in L$ and $t \in Z^l(L)$ as follows:

$$\begin{aligned} & [x_{\lambda_1} l_{\lambda_2, \lambda_3}^3(y, z, t)] - [y_{\lambda_2} l_{\lambda_1, \lambda_3}^3(x, z, t)] \\ & + [z_{\lambda_3} l_{\lambda_1, \lambda_2}^3(x, y, t)] + [l_{\lambda_1, \lambda_2}^3(x, y, z)_{\lambda_1+\lambda_2+\lambda_3} t] - l_{\lambda_1+\lambda_2, \lambda_3}^3([x_{\lambda_1} y], z, t) \\ & - l_{\lambda_2, \lambda_1+\lambda_3}^3(y, [x_{\lambda_1} z], t) - l_{\lambda_2, \lambda_3}^3(y, z, [x_{\lambda_1} t]) \\ & + l_{\lambda_1, \lambda_2+\lambda_3}^3(x, [y_{\lambda_2} z], t) + l_{\lambda_1, \lambda_3}^3(x, z, [y_{\lambda_2} t]) - l_{\lambda_1, \lambda_2}^3(x, y, [z_{\lambda_3} t]) \end{aligned}$$

$$\begin{aligned}
&= -\frac{1}{4} \left(\llbracket x_{\lambda_1} ((y \circ z) \circ_{\lambda_2+\lambda_3} t) \rrbracket - \llbracket y_{\lambda_2} ((x \circ_{\lambda_1} z) \circ_{\lambda_1+\lambda_3} t) \rrbracket \right. \\
&\quad + \llbracket z_{\lambda_3} ((x \circ_{\lambda_1} y) \circ_{\lambda_1+\lambda_2} t) \rrbracket - 0 - (\llbracket x_{\lambda_1} y \rrbracket \circ_{\lambda_1+\lambda_2} z) \circ_{\lambda_1+\lambda_2+\lambda_3} t \\
&\quad + (\llbracket x_{\lambda_1} z \rrbracket \circ_{\lambda_1+\lambda_3} y) \circ_{\lambda_1+\lambda_2+\lambda_3} t - (y \circ_{\lambda_2} z) \circ_{\lambda_2+\lambda_3} \llbracket x_{\lambda_1} t \rrbracket \\
&\quad - (x \circ_{\lambda_1} \llbracket y_{\lambda_2} z \rrbracket) \circ_{\lambda_1+\lambda_2+\lambda_3} t + (x \circ_{\lambda_1} z) \circ_{\lambda_1+\lambda_3} \llbracket y_{\lambda_2} t \rrbracket \\
&\quad \left. - (x \circ_{\lambda_1} y) \circ_{\lambda_1+\lambda_2} \llbracket z_{\lambda_3} t \rrbracket \right) \\
&= -\frac{3}{8} (x \circ_{\lambda_1} (y \circ_{\lambda_2} z) - (x \circ_{\lambda_1} y) \circ_{\lambda_1+\lambda_2} z - y \circ_{\lambda_2} (x \circ_{\lambda_1} z)) \circ_{\lambda_1+\lambda_2+\lambda_3} t \\
&= 0.
\end{aligned}$$

The other cases can be checked similarly. The proof is completed. $\square$

It is natural to introduce the general notions of L_∞ -conformal algebras and $Leib_\infty$ -conformal algebras. In fact, an initial version of this paper appeared in arXiv: [2301.06903](#). Since then, several articles developing the theory presented in this manuscript have been appeared, such as those in [AW24; DS23; SD23]. We encourage interested readers to consult these references for further developments in the field.

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